

# The Monte Carlo Independent Column Approximation: An assessment using several global atmospheric models

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**ABSTRACT:** The Monte Carlo Independent Column Approximation (McICA) computes domain-average, broadband radiative flux profiles within conventional global climate models (GCMs). While McICA is unbiased with respect to the full ICA, it generates, as a by-product, random noise. If this by-product leads to statistically significant impacts on GCM simulations, it could limit the usefulness of McICA. This paper assesses the impact of McICA's random noise on six GCMs. To this end, the GCMs performed ensembles of 14-day long simulations for various renditions of McICA, each with differing amounts of random noise. As seen in the past, low-cloud fraction and surface temperature were affected most by noise. However, all GCM simulations using operationally viable renditions of McICA showed no statistically significant impacts, even for precipitation – a highly intermittent variable that one might expect to be sensitive to random fluctuations. Two GCMs showed statistically significant responses using an *academic* version of McICA that generates overly large sampling noise. Time series analyses of high-resolution (i.e. typically 2-hourly) data revealed that fluctuations associated with most variables and GCMs are immune to McICA noise. Moreover, the nature of these fluctuations can vary substantially among GCMs and most often they overwhelm any noise impacts. Overall, the results presented here corroborate a range of previous studies done on one GCM at a time: random noise produced by recommended versions of McICA has statistically insignificant effects on GCM simulations. Copyright © 2008 Royal Meteorological Society and Her Majesty in Right of Canada.

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## 1. Introduction

### 1.1. Background

It is recognized widely that, if general circulation models (GCMs) are to satisfactorily simulate Earth's climate, physical processes responsible for hydro-radiative feedback mechanisms must be addressed and represented well (e.g. Aires and Rossow, 2003; Randall *et al.*, 2003). Failing to cross this ill-defined line, a GCM stands a good chance of serving only to heighten confusion surrounding explanation, and thus simulation, of the Earth–atmosphere system (cf. IPCC, 2007). Central to adequate representation of hydro-radiative feedbacks lies numerical simulation of atmospheric radiative transfer (e.g. Stephens, 2005).

The common purpose for doing radiative transfer calculations in GCMs is to provide atmospheric and surface broadband flux convergence rates. Since the mid-1970s, GCMs have employed multi-layer two-stream approximations (TSAs) of the radiative transfer equation (RTE; Meador and Weaver, 1980). At the heart of TSAs is the assumption that atmospheric layers are horizontally homogeneous. Hence, radiation flowing within a GCM column is assumed not to interact with neighbouring columns. This is the independent column approximation (ICA; Stephens *et al.*, 1991). Given that horizontal grid spacings exceed 50 km in most GCMs, it is reasonable to apply the ICA at a GCM's inner scale.

Given the task at hand, multi-layer TSA algorithms eventually came to address columns whose layers were only partially filled by clouds which, in turn, overlapped vertically according to idealized rules, notably maximum-random overlap (MRO; Geleyn and Hollingsworth, 1979; Morcrette and Fouquart, 1986; Tian and Curry, 1989; Stubenrauch *et al.*, 1997). Given that MRO systematically underestimates vertically projected

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cloud fraction (Räisänen *et al.*, 2004), GCM radiative transfer codes based on MRO *rely* on horizontally homogenous clouds that are too attenuative relative to variable clouds with the same mean optical thickness (cf. Barker *et al.*, 1999, 2003; Cole *et al.*, 2005).

While several attempts have been made to account for effects of horizontal variations in single-layer clouds (e.g. Stephens, 1988; Davis *et al.*, 1990; Cahalan *et al.*, 1994; Barker, 1996; Cairns *et al.*, 2000), they resemble each other in that they use the TSA, strive for analytic expressions for layer reflectance and transmittance, but pay little attention to covariations in the horizontal and vertical. Oreopoulos and Barker (1999) recognized the untenability of investing the MRO scheme with horizontally variable clouds, but their model and results resembled other attempts to account for unresolved cloud fluctuations in 1D codes: limited ranges of applicability and, seemingly, irreconcilable biases relative to the ICA.

Judging from GCM radiative transfer algorithms developed from the late 1970s to early 2000s, it could be argued that a paradigm was established based on the logic:

Because radiative fluxes depend on cloud structure, this implies that descriptions of cloud structure at scales  $\Delta x < \Delta x_{\text{GCM}}$ , where  $\Delta x_{\text{GCM}}$  is the GCM grid spacing, should be a subset of (i.e. built directly into) the radiative transfer algorithm.

There is, however, no sound reason to believe that this logic will lead to algorithms that either allow for satisfactory description of cloud properties or account properly for interactions between radiation and clouds at scales smaller than  $\Delta x_{\text{GCM}}$ . Indeed, models based on this logic seem bound to approximations, limitations, and biases.

## 1.2. The Monte Carlo-based solution

If a GCM's column were to be resolved into  $J$  sub-columns, and the ICA applied to them, domain-average, broadband fluxes would be defined as

$$\langle F \rangle = \frac{1}{J} \sum_{j=1}^J \sum_{k=1}^K F(j, k), \quad (1)$$

where  $K$  is the number of spectral intervals, and  $F$  is flux from a 1D solution of the RTE. While (1) is the standard aspired to by GCM radiation algorithms (Collins, 2001), it is computationally intractable. As an alternative to (1), and the aforementioned 1D solvers, Barker *et al.* (2002) proposed the Monte Carlo ICA (McICA) method. McICA rests on sampling stochastically generated subcolumns of a GCM column, performing 1D radiative transfer calculations on them, and averaging the results (Pincus *et al.*, 2003, 2006; Räisänen *et al.*, 2004). As sampling is done during spectral integration, McICA reduces (1) to a single sum over  $k$ .

McICA represents a distinct departure from conventional methods because it separates the description of

unresolved media from the radiative transfer algorithm while remaining unbiased with respect to (1). As such, it is extremely flexible for it admits any one-point statistical description of unresolved media, and subgrid variability can be handled easily by any 1D solution of the RTE and not just TSAs. Moreover, because a basic TSA (e.g. Wiscombe, 1977) is used, McICA execution times are close to, if not less than, those of conventional schemes. As a by-product, however, McICA produces flux profiles that contain sampling noise.

Three studies (Räisänen *et al.*, 2005; Pincus *et al.*, 2006; Morcrette *et al.*, 2008) suggest that McICA noise can be reduced to the point of having negligible impact on GCM simulations. Räisänen *et al.*'s (2005) tests with the National Center for Atmospheric Research (NCAR) Community Atmospheric Model version 1.8 (CAM-v1.8) showed however that, if large enough, McICA's noise can have statistically significant affects on GCM simulations. Thus, the main point of the current study is to assess the sensitivities of several GCMs, with differing physical parametrizations and dynamical cores, to stochastic noise generated by McICA. To this end, each GCM, with various renditions of McICA and a suitable stochastic subgrid-scale cloud generator, performed 14-day ensemble simulations.

The following two sections describe the experiments and the GCMs involved in the intercomparison. The fourth section presents the magnitudes of McICA noise and summarizes the results of the experiments, and the final section contains conclusions.

## 2. Experimental design

Table 1 summarizes the GCMs used in this study. The experiments performed by each GCM follow those performed by Räisänen *et al.* (2005). These experiments are summarized in the following subsections, but first the rationale for the experiments is discussed.

### 2.1. Rationale

Stochastic fluctuations of radiative fluxes generated by McICA can be thought of as extraneous information injected into a GCM at its radiative time step and grid spacing. For a fairly noisy version of McICA, Räisänen and Barker (2004) estimated typical standard deviations for instantaneous net surface irradiances to be  $\sim 50 \text{ W m}^{-2}$ , the vast majority due to solar radiation. From time step to time step and cell to cell, McICA noise is uncorrelated so standard deviations of accumulated fluxes decay like  $1/\sqrt{N}$ , where  $N$  is number of time steps or grid points. Given that the radiative time step for most GCMs is  $\sim 1 \text{ h}$ , the magnitude of McICA noise for diurnal-mean net surface fluxes is  $O(10 \text{ W m}^{-2})$ . As can be seen in Figure 1, this is much smaller than diurnal standard deviations for surface irradiance that arise via the diurnal solar cycle (which are approximately half the values shown). It stands to reason, therefore, that for McICA's noise to excite change, it would have to

Table I. Summary of the characteristics of participating GCMs.

Institute	Model	Resolution	Dynamic time step (min)	Save interval (hr)
Environment Canada–Canadian Meteorological Centre (EC-CMC)	GEM	$1.5^{\circ} \times 1.5^{\circ}$ ; L60	30	2
Environment Canada–Canadian Centre for Climate Modelling and Analysis (EC-CCCma)	CCCma	T47; L35	20	2
Finnish Meteorological Institute (FMI)	ECHAM5-FMI	T42; L31	20	2
European Centre for Medium-range Weather Forecasts (ECMWF)	ECMWF	$T_L 159$ ; L91	60	2
Geophysical Fluid Dynamics Laboratory (GFDL)	GFDL-AM2	$2^{\circ} \times 2.5^{\circ}$ ; L24	60	1 and 3
Pennsylvania State University (PSU)	CAM3_PSU	T42; L26	60	1

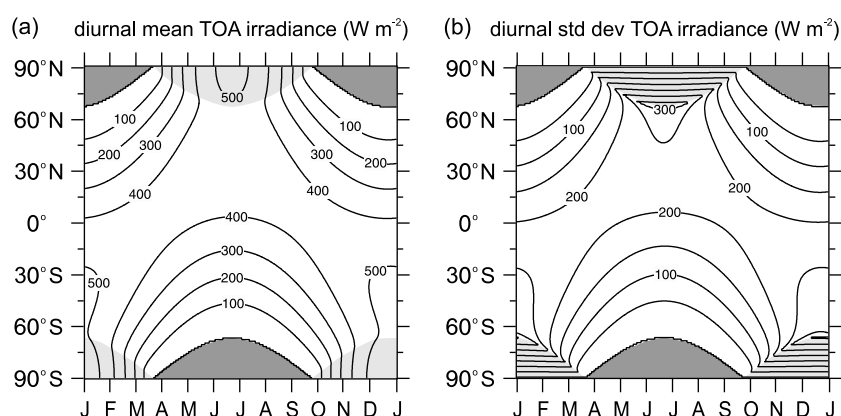


Figure 1. (a) diurnal-mean solar irradiance at the top of the atmosphere (TOA) as a function of time and latitude. (b) shows the corresponding standard deviation of solar irradiance at the TOA due simply to the Earth's rotation. Corresponding values for net surface irradiance are about half those shown here. Light and dark shaded regions denote perpetual sun-up and sun-down periods, respectively.

do so quickly. Likewise, if noise-induced, short-term impacts yield significant biases, one can expect those biases to effectively force slower climatic variables and thus jeopardize McICA's utility.

Following this line of reasoning, relatively short simulations should suffice to elucidate impacts of McICA's stochastic noise on *fast* components of the climate system. Since *slower* components, such as deep soil moisture and mixed-layer ocean temperature, respond to surface energy budgets integrated over weeks, it is hypothesized that if fast components do not respond to high-frequency noise, long-term climate simulations will not respond either. This rationale resembles that of the US DoE's Climate Change Prediction Program (CCPP) and Atmospheric Radiation Measurement (ARM) Program (CCPP-ARM) Parameterization Testbed (CAPT) study (Philips *et al.*, 2004).

## 2.2. The experiments

Based on the reasoning just explained, 14-day long GCM simulations were used for this intercomparison. McICA fluctuations for surface fluxes integrated over two weeks are about 20 times smaller than instantaneous values. Following Räisänen *et al.*'s (2004) calculations, this corresponds roughly to radiative fluctuations associated with fortnightly mean total cloud fraction changes of  $\sim 0.02$ .

Hence, we are satisfied that if significant differences do not appear after 14 days, it is unlikely that the noise could somehow tunnel-up at longer scales and significantly alter a simulation relative to a noiseless reference.

There are several versions of McICA that yield substantially different amounts of noise (Räisänen and Barker, 2004). The least noisy incarnation was defined as the reference. Most GCMs performed five experiments with varying amounts of noise, each consisting of a 10-member ensemble. The experiments are described in the following subsections. Ensembles were created for each experiment by starting the GCM 6 hours apart off a 'spun-up' simulation. As such, the  $j$ th ensemble member of all experiments had the same initial conditions. This initialization process is illustrated in Figure 2.

Two sets of experiments were performed: one started from 1 January, the other from 1 July. This was to explore whether the impact of McICA stochastic noise is state-dependent in addition to GCM-dependent. All indications suggest that this is not the case so only January results are shown and discussed. In all cases, sea-surface temperatures were prescribed. For brevity, the following explanation of experiments assumes a radiation code based on the correlated  $k$ -distribution (CKD) method (Fu and Liou, 1992), which was not always the case.

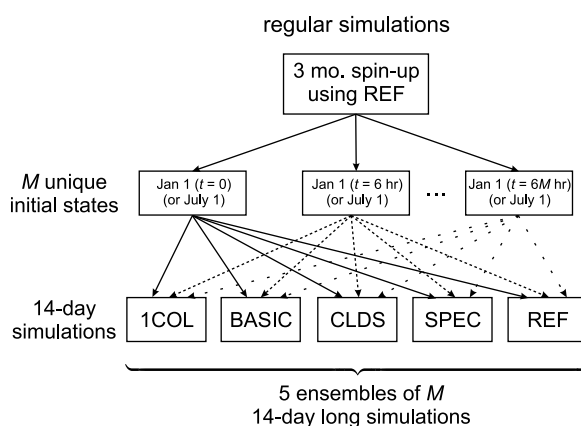


Figure 2. Schematic diagram showing production of ensembles of 14-day simulations. All GCMs produced 10-member ensembles (i.e.  $M = 10$ ) though not all performed all five ensembles which are denoted here as 1COL, BASIC, CLDS, SPEC and REF.

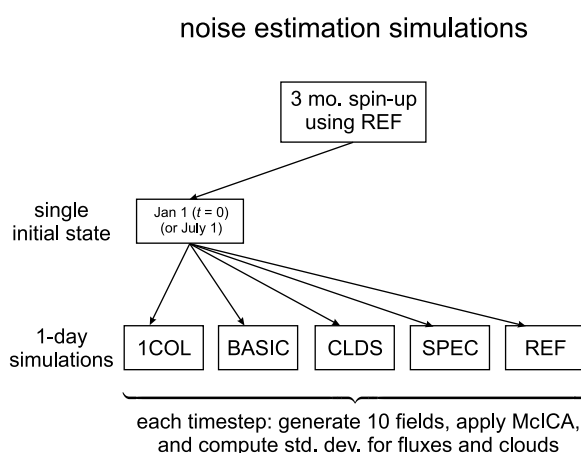


Figure 3. Schematic diagram showing production of noise estimates for various renditions of McICA (denoted here as 1COL, BASIC, CLDS, SPEC and REF).

### 2.2.1. Experiment 1: 1COL

If a *single* subcolumn is generated stochastically at each time step and spectral integration (over  $K$  quadrature points) is performed on that subcolumn, (1) becomes

$$\langle F \rangle = \sum_{k=1}^K F(\alpha, k) \quad (2)$$

where  $\alpha$  represents *the* generated subcolumn. This approach yields copious McICA noise and is referred to as '1COL'. Moreover, it resembles the method used in the Goddard Institute for Space Studies GCM (A. A. Lacis, 2007, personal communication).

### 2.2.2. Experiment 2: BASIC

Sampling is improved over 1COL by generating a subcolumn (cloudy or clear) for each integration point and computing flux as

$$\langle F \rangle = \sum_{k=1}^K F(\alpha_k, k) \quad (3)$$

where  $\alpha_k$  represents one of  $K$  generated columns. Being the most straightforward, code-friendly version of McICA, it is referred to as the 'BASIC' experiment.

### 2.2.3. Experiment 3: CLDS

Because GCMs have to compute clear-sky radiative fluxes to estimate cloud radiative effects, it is reasonable, though not necessary, to perform a full clear-sky calculation and devote all of McICA's samples to the cloudy (variable) portion of a domain. Thus, by demanding that randomly generated subcolumns contain cloud, total flux is

$$\langle F \rangle = (1 - A_c) \sum_{k=1}^K F(\alpha_{\text{clr}}, k) + A_c \sum_{k=1}^K F(\alpha_{\text{cld},k}, k), \quad (4)$$

where  $\alpha_{\text{clr}}$  represents *the* clear-sky subcolumn,  $\alpha_{\text{cld},k}$  represents one of  $K$  cloudy subcolumns, and  $A_c$  is vertically projected cloud fraction.  $A_c$  can be estimated rapidly and accurately by generating many subcolumns with a stochastic cloud generator that avoids estimation of water contents. Since random sampling is confined to cloudy subcolumns, this experiment is referred to as CLDS. Obviously, the amount of noise generated by BASIC is greater than or equal to that produced by CLDS.

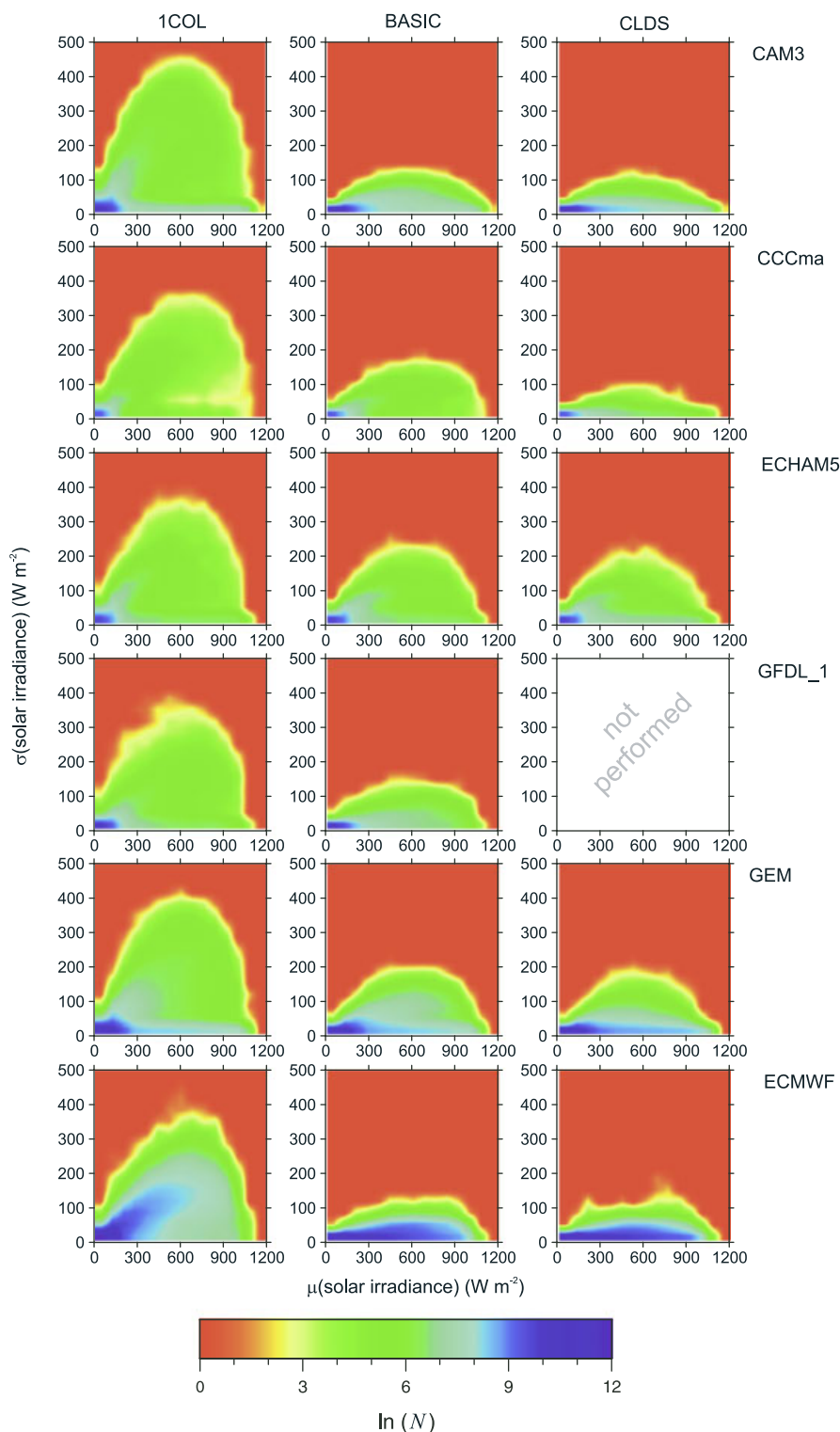


Figure 4. Standard deviation  $\sigma$  of surface solar irradiance as a function of mean solar irradiance  $\mu$  for three versions of McICA (1COL, BASIC and CLDS) for the six GCMs considered in this study. Colours indicate the natural logarithm of the number of samples in each bin (of which there are  $15 \times 15$ ). This figure is available in colour online at [www.interscience.wiley.com/journal/qj](http://www.interscience.wiley.com/journal/qj)

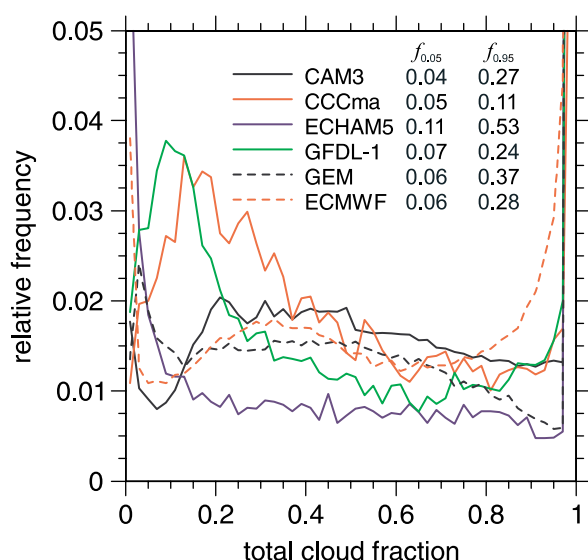


Figure 5. Area-weighted relative frequencies of total cloud fraction for the entire Earth for the six GCMs used in this study.  $f_{0.05}$  and  $f_{0.95}$  refer to the fractional areas with total cloud fractions less than 0.05 and greater than 0.95. This figure is available in colour online at [www.interscience.wiley.com/journal/qj](http://www.interscience.wiley.com/journal/qj)

#### 2.2.4. Experiment 4: SPEC

An efficient way to reduce McICA noise is to sample additional cloudy subcolumns for spectral intervals that contribute most to McICA's noise. These points are generally those with large cloud radiative effects. As such, this version of McICA defines fluxes as

$$\langle F \rangle = (1 - A_c) \sum_{k=1}^K F(\alpha_{\text{clr}}, k) + A_c \sum_{k=1}^K \left[ \frac{1}{N_k} \sum_{i=1}^{N_k} F(\alpha_{\text{cld}, i, k}, k) \right], \quad (5)$$

where  $N_k$  are the number of samples generated for  $k$ , and  $\mathcal{N} = \sum_{k=1}^K N_k$  is the total number of samples generated. This version of McICA is referred to as SPEC because additional samples are assigned to specific spectral intervals (cf. the splitting technique as described in Marchuk *et al.*, 1980). Note that for CLDS,  $N_k = 1$  and so  $\mathcal{N} = K$ . However, it was recommended that for this model  $\mathcal{N} \approx 1.5K$ . Räisänen and Barker (2004) provide a procedure for defining  $\{N_k\}$  that reduces variance optimally.

#### 2.2.5. Experiment 5: REF

Finally, a reference experiment, referred to as REF, was performed by most GCMs, and characterized by very weak noise. The simplest and most economical way to do this is to use the SPEC model with large  $\mathcal{N}$ . The GFDL GCM uses a very small number of spectral points and therefore its REF experiment was the (less economical) ICA using  $K$  subcolumns (Pincus *et al.*, 2006). Nevertheless, noise associated with REF is much smaller than that of 1COL and notably less than that of SPEC.

### 3. Description of GCMs

Tables I and II summarize characteristics of the GCMs that participated in the intercomparison. CMC and GFDL are finite-volume models; the other four are spectral models. Horizontal resolutions varied from  $\sim 100$  km to  $\sim 250$  km, and number of vertical layers ranged from 24 to 91. While some subgrid-scale parametrizations are common to more than one GCM (e.g. GEM and CCCma used the same radiation codes), and others use parametrizations with common lineages (e.g. ECMWF and ECHAM5), no two GCMs share the exact same collection of parametrizations.

Table II. Summary of the characteristics of GCM radiative transfer models. (Radiation time step = 1 h in all cases.)

Model	Quadrature points (SW/LW)
GEM	35/ 46
CCCma	35/ 46
ECHAM5-FMI	4/140
ECMWF	112/140
GFDL	18/ 7
CAM3_PSU	55/ 67

Morcrette (2000) and Räisänen *et al.* (2005) showed that changing the radiative time step  $\Delta t_{\text{rad}}$  from 1 h to 3 h can influence a GCM's simulation. With McICA there is the added concern that large  $\Delta t_{\text{rad}}$  might allow a GCM to incorporate the effects of radiative anomalies into the overall character of the simulation (Räisänen *et al.*, 2005). For the experiments reported here, the GFDL and GEM models were run with 3 h and 2 h time steps in addition to a 1 h time step. Differences between the runs were very minor, so results only for the shorter time step are discussed below.

Lacking rules pertaining to description of unresolved horizontal variations in cloud water and vertical overlap of fractional cloud, groups were left to define them for themselves. While this sidesteps the utility of McICA, recall that it is the *spread* in McICA-generated stochastic noise that is important for these experiments; details of unresolved cloud morphology are secondary. As seen later, wide ranges of noise were achieved by all GCMs.

### 4. Results

Results are presented in two sections. The first summarizes variations in stochastic noise as a function of GCM. The second presents GCM responses as a function of noise.

#### 4.1. Magnitude of McICA noise

To quantify the magnitude of the stochastic noise generated for the five experiments, standard deviations of variables responsible for producing noise were computed for each rendition of McICA by generating subgrid-scale



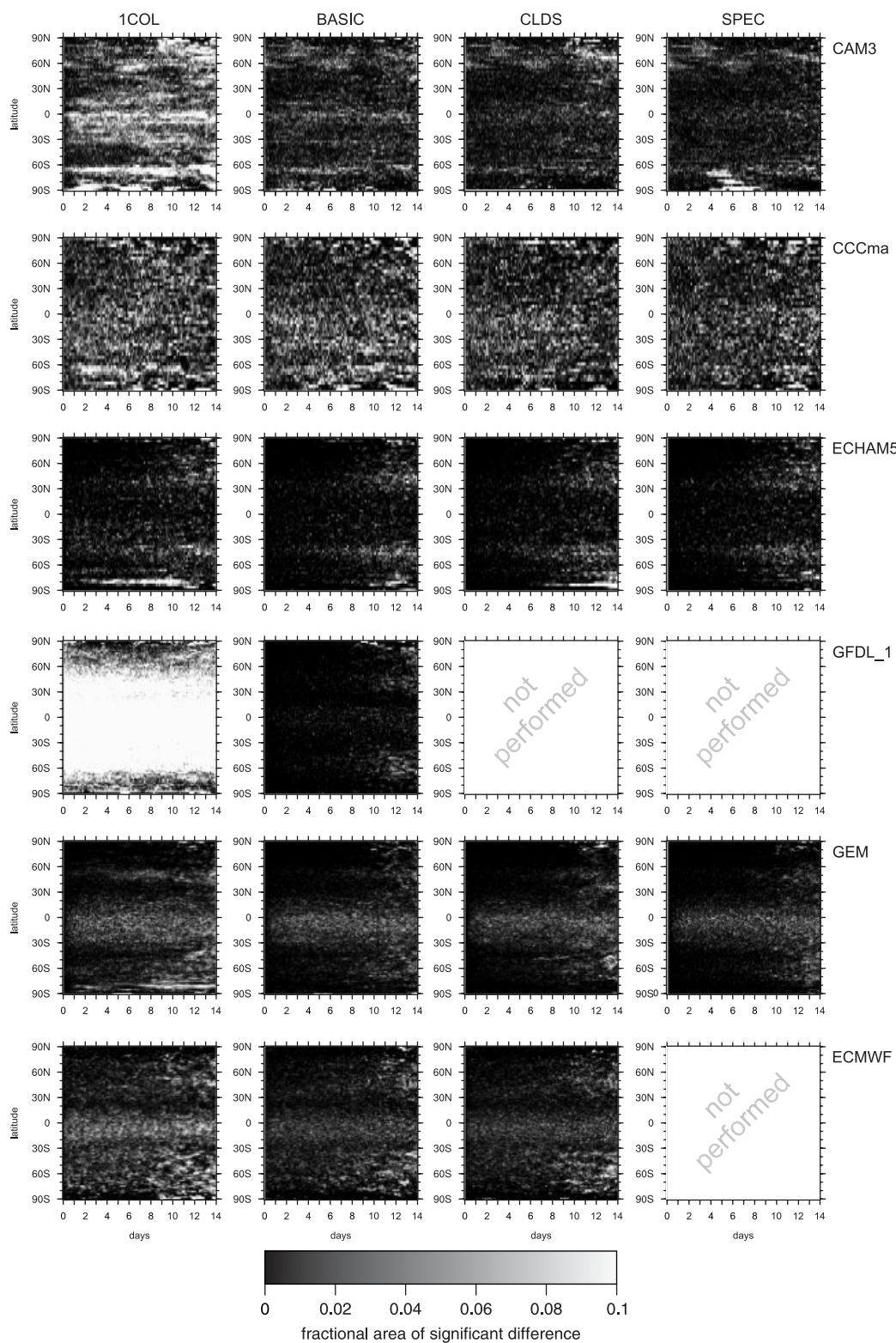


Figure 6. Hovmöller-like diagrams of fractional areas of zonal bands exhibiting statistically significant differences, at the 95% confidence level, in low-cloud fraction for the various renditions of McICA relative to their respective reference (REF) simulations. McICA noise decreases from left to right. Note that if samples were drawn only from the REF simulation (i.e. the experimental samples were unambiguously pulled from the control population), plots would be characterized by white noise, resembling 'snow' on a TV screen, with means of 0.05.

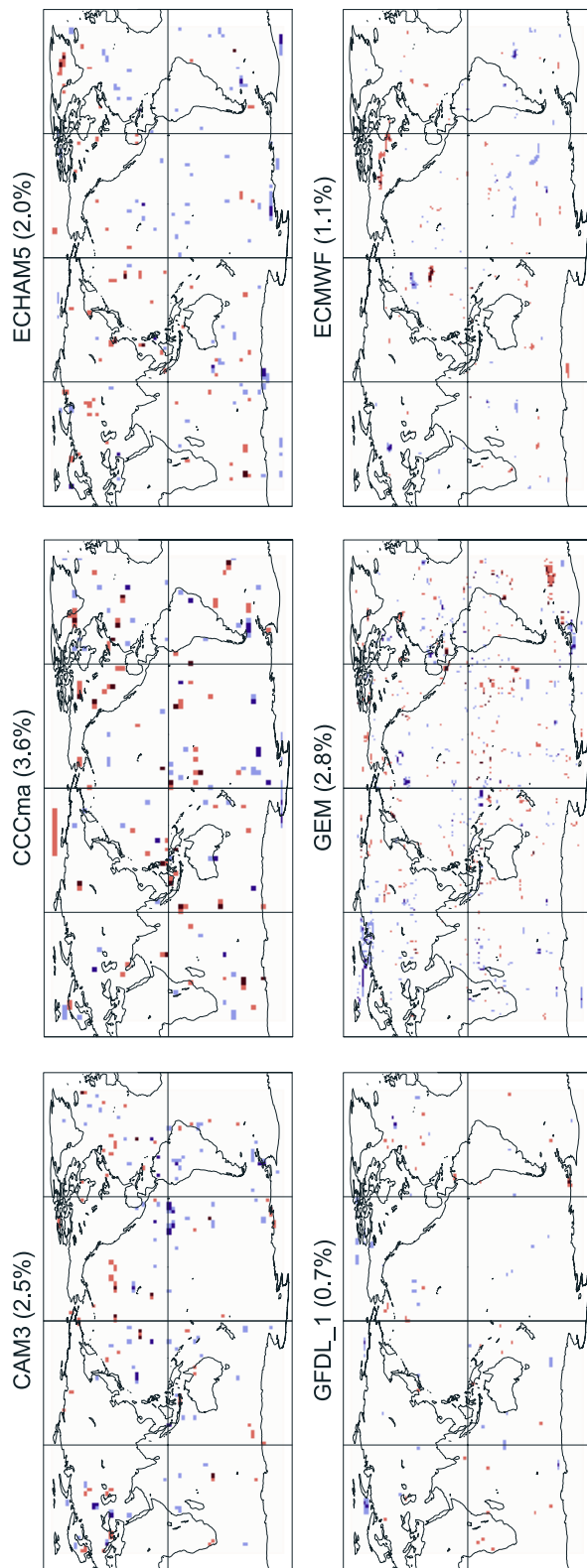


Figure 7. Global plots showing statistically significant differences between mean low-cloud fractions on the 14th day of the simulations for the BASiC rendition of McICA relative to the corresponding REF control. Light blue and red correspond to under- and overestimations by BASiC that are significant at the 95% confidence level. Likewise, dark blue and red signify the same but at the 99% confidence level. The GCM is indicated in the title along with the fractional area of the globe having statistically significant differences at the 95% confidence level. This figure is available in colour online at [www.interscience.wiley.com/journal/qj](http://www.interscience.wiley.com/journal/qj)



clouds and calling the radiation code ten times at each radiative time step for a *single day* (Figure 3). The variables are short-wave (SW) and long-wave (LW) radiative flux profiles, total cloud fraction, and cloud liquid and ice water paths (actually, water contents are subject to noise directly).

As an illustrative example, Figure 4 shows standard deviation of SW radiative flux at the surface as a function of the corresponding mean for the 1COL, BASIC, and CLDS versions. The magnitude of noise and the rate at which it decreases with increasing samples (left to right across the figure) depends on several factors including number of spectral intervals, diurnal cycle of cloud properties, frequency distribution of cloud fraction profile, and assumptions about unresolved clouds. This plot shows the abundance of noise associated with 1COL as well as the dramatic, and ubiquitous, reduction of noise when going from 1COL to BASIC. The only GCM to show a sizable reduction in noise when going from BASIC to CLDS, where the total number of samples remained the same, is the CCCma GCM. The most likely reason for this has to do with the distribution of total cloud fraction; McICA errors maximize at intermediate cloud fractions and vanish at clear sky and almost vanish at overcast. As Figure 5 shows, the CCCma GCM has the fewest near-overcast and cloudless cases and a large fraction of total cloud fractions near 0.25. The other GCMs have relatively few intermediate total cloud fractions; especially ECHAM5 which produces numerous near-clear and overcast events. Note that a GCM that produces *only* clear skies and overcasts will realize minimal variance reduction when going from BASIC to CLDS. This situation should become increasingly prevalent as  $\Delta x_{\text{GCM}} \rightarrow 0$ .

From Figure 4 alone one can predict, with reasonable confidence, that most GCMs have little to gain from CLDS, let alone the more elaborate SPEC. Nevertheless, as argued by Räisänen and Barker (2004), all GCMs compute clear-sky fluxes for diagnostic reasons, so there is some motivation to employ CLDS and devote all available stochastic samples to the cloudy, variance-generating, portion of a column.

#### 4.2. Impacts on GCM simulations

While numerous fields were saved from the simulations, so few exhibited significant differences that presentation of results focuses on low-cloud fraction and 2 m air temperature, the variables that showed the greatest sensitivities to noise. Results are presented in two subsections: one focusing on fractional areas of the globe that exhibit statistically significant changes due to noise, and the other on time series analysis.

##### 4.2.1. Fractional areas of statistically significant differences

A convenient way to summarize the kind of results obtained here is the paired-difference version of Student's *t*-test (e.g. von Storch and Zwiers, 1999). This is because

ensemble members in one experiment have the same initial conditions as those in another experiment. Hence, the experiments are not fully independent of each other. If two experiments, each with  $N$  ensemble members, are being compared and  $\mathbf{x}_n$  and  $\mathbf{y}_n$  are the  $n$ th member of their respective ensembles, define  $\mathbf{d}_n = \mathbf{x}_n - \mathbf{y}_n$ . The null hypothesis to be tested is therefore  $H_0: \mu_{\mathbf{xy}} = 0$ , where  $\mu_{\mathbf{xy}}$  is the mean of the population of differences  $\mathbf{d}$  from which the sample, of size  $N$ , was drawn. From sampling theory, the optimal test statistic is

$$t = \frac{\hat{\mu}_{\mathbf{xy}}}{S_{\mathbf{xy}}/\sqrt{N}}, \quad (6)$$

where  $\hat{\mu}_{\mathbf{xy}}$  is the sample mean difference, and  $S_{\mathbf{xy}}$  is the corresponding unbiased sample variance. Assuming that the sample differences follow Student's distribution,  $t$  is checked to see if it lies beyond a specified value, the significance level, in the tail of the distribution. This is done by checking if

$$\frac{1}{v^{\frac{1}{2}} B\left(\frac{1}{2}, \frac{v}{2}\right)} \int_{-t}^t \left(1 + \frac{x^2}{v}\right)^{-\frac{v+1}{2}} dx > t_{\text{crit}}, \quad (7)$$

where  $v = N - 1$ ,  $B(1/2, v/2)$  is the beta function, and  $t_{\text{crit}} = 1 - \alpha$ , where  $\alpha$  is the level of significance one is testing for. If (7) is true, the sampled difference is considered to be statistically significant. Throughout this study,  $\alpha = 0.05$  implying that all reported tests are at the 95% confidence level. Note that if  $M$  samples are drawn at random from a population, and sample pairs are tested via (6) and (7), then on average,  $\alpha M$  of them will exhibit statistically significant differences. Reported here are fractional areas of *regions* with GCM cells that exhibit statistically significant differences. The *regions* are typically zonal bands and the entire Earth.

Räisänen *et al.* (2005) showed that for CAM1.8, low-cloud fraction exhibited the greatest sensitivity to McICA noise. Low cloud is defined as cloud below 680 hPa. Figure 6 shows Hovmöller-like diagrams of the fractional area of zonal bands with statistically significant differences in low-cloud fraction for four renditions of McICA relative to REF. The striking aspect of this set of plots is that each model responds differently to McICA noise, though all but some of the 1COL versions display statistically insignificant ramifications of noise. For instance, CCCma and ECHAM5 show almost no dependence on noise level, while CAM3 and GFDL experience dramatically reduced sensitivities to noise going from 1COL to BASIC. The two highest-resolution models, GEM and ECMWF forecast models, show only a slight dependence on noise level, but, despite being very different models, they share some common features – namely, remarkably few significant responses in the subtropical high regions with persistent, though minor, differences throughout the Tropics. Presumably differences in the Tropics stem from local processes responding very quickly to noise but of such limited strength that their influences are not felt very far outside the cell in which they were initiated.

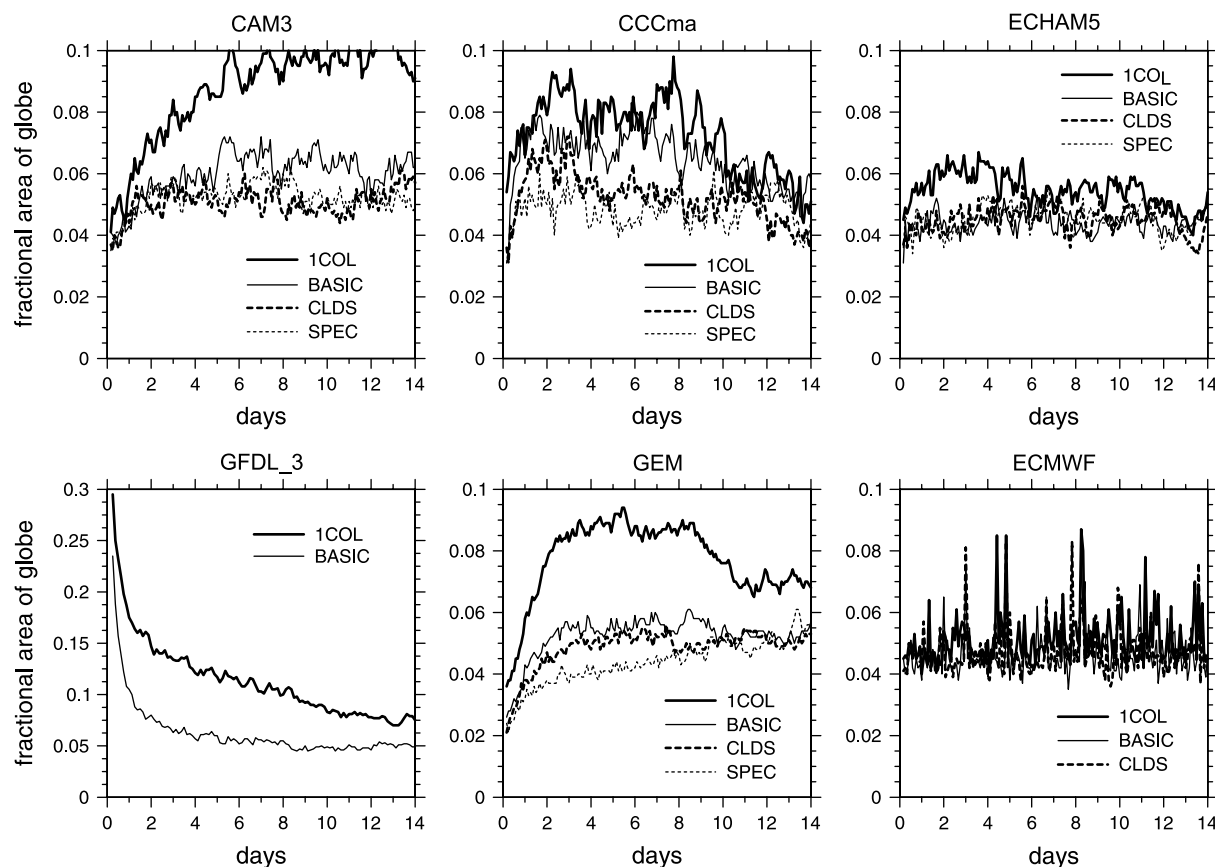


Figure 8. Fractional areas of the globe exhibiting statistically significant 2 m air temperature differences, at the 95% confidence level, relative to the REF control as functions of time for all experiments performed by all GCMs. Had the experiments been drawn from the control population, lines would have been noisy with means very close to 0.05 (as most of the curves are beyond about the 10th day).

Several models (ECMWF, ECHAM5, GEM and GFDL) show similar patterns for the mid and high latitudes: a rather sudden transition, at roughly 8 days into the simulations, from areas of difference slowly increasing with time to being independent of time. Nevertheless, most latitudes have areas of significant difference of approximately 5%, so these appear to be fluctuations due simply to sampling from the same population with little, or nothing, to do with response to noise.

The results just shown are applicable to 2 h instantaneous samples. Figure 7 shows global maps of the location of significant differences in low-cloud fraction averaged over the 14th day of the BASIC simulations. In general, the fractional areas showing significant differences relative to REF are much smaller than at the 2-hourly zonal level; the largest area is a meagre 3.6% for the CCCma model. Moreover, one is hard-pressed to discern any pattern in any of these plots. This is another positive feature, for it suggests that the effects of noise have not been incorporated and transferred up-scale to affect change on regional scales.

Figure 8 shows time series of fractional areas of the globe that exhibit statistically significant noise-induced responses on 2 m air temperature as functions of time. As with low cloud, each model displays a unique response to noise. Most models and renditions of McICA begin

the simulation looking as though they were drawn from the same population as REF. The 1COL versions of CAM3 and GEM show slight impacts of noise, while the corresponding versions of CCCma and GFDL appear to be heading toward acceptance of  $H_0$  by day 14. Consistent with Räisänen *et al.*'s (2005) analyses, all other cases are statistically indistinguishable from REF.

Precipitation is another crucial climatic variable that one would hope is not influenced by noise. Moreover, precipitation fields differ vastly from cloud fraction and temperature on account of their extreme intermittency. Figure 9 shows Hovmöller-like diagrams of fractional areas of significant difference relative to REF for large-scale precipitation for three GCMs and total precipitation for the remaining three. Only 1COL and BASIC are shown. Plotted along the bottom are time series of corresponding global areas. Most plots and curves demonstrate extremely minor affects due to noise, often even smaller than impacts on 2 m temperature. The exception is ECMWF's run with 1COL which shows a marked response in the Tropics. Again, however, it is highly unlikely that a modelling group would use 1COL, and the impact is clearly crushed by the noise reduction of BASIC, which is ECMWF's operational configuration (Morcrette *et al.*, 2008).

As a final means of portraying the ramification, or lack thereof, of McICA noise, Figure 10 shows fractional areas

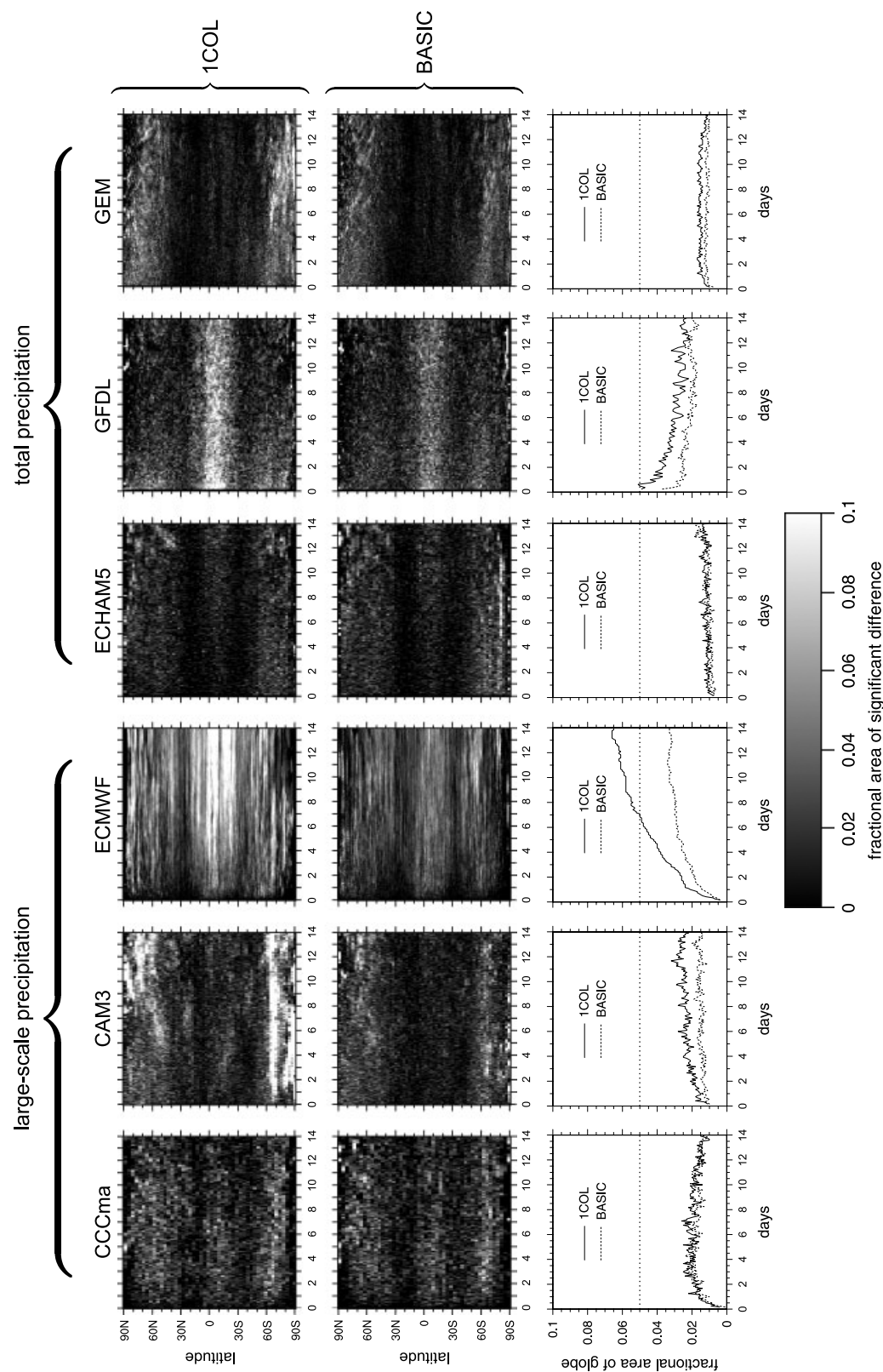


Figure 9. Along the top are two rows of Hovmöller-like diagrams of fractional areas of zonal bands exhibiting statistically significant differences, at the 95% confidence level, in precipitation for the 1COL and BASIC renditions of McICA relative to their respective reference (REF) simulations. Had samples been drawn only from the REF simulation, plots would have been characterized by white noise with a mean of 0.05 (corresponding to the 95% confidence level). The lower row of line plots show corresponding fractional areas of the globe exhibiting statistically significant precipitation differences, at the 95% confidence level, relative to the REF control as functions of time. Had the experiments been drawn from the control population, lines would have been noisy with means very close to 0.05.

of the globe showing statistically significant changes to low cloud fraction and 2 m temperature averaged over the 14th day as functions of globally averaged McICA noise associated with net short-wave flux at the surface. These plots show succinctly that if ever there is an issue with noise, it is bound to the 1COL version of McICA. They also show that the CAM stands the most to gain by going from 1COL to BASIC, while the CCCma stands the most to gain by going from BASIC to CLDS.

#### 4.2.2. Time series analysis

An obvious area to look for rapid development of responses due to high-frequency noise generated by McICA is in the magnitude of temporal fluctuations for variables  $\varphi$  with short characteristic response times. Variables in the lower atmosphere or at the surface are good candidates, particularly ones that are influenced directly by McICA-related noise. For a discrete time series spanning a length of time  $L$ , fluctuations of  $\varphi$  are assessed via the second-order structure function which is defined as

$$S_2(r) = \langle [\varphi(t+r) - \varphi(t)]^2 \rangle, \quad (8)$$

where  $t$  is time and  $r$  is a time lag.  $S_2(r)$  is related, by the Wiener–Khinchine theorem, to the power spectrum (Marshak and Davis, 2005), and is a very intuitive measure of typical fluctuations over a wide range of time. For ranges of  $r$  that behave like

$$S_2(r) \sim r^{\zeta(2)}, \quad (9)$$

$1 < \beta = \zeta(2) + 1 < 3$ , where  $\beta$  is the spectral exponent of the power spectrum.

The region of central Canada bounded by 55°–65°N and 95°–105°W has a continental winter climate with surface temperatures that are sensitive to the presence of cloud. Averaging over this relatively large area, which is about seven times as large as the coarsest GCM grid spacing, should alleviate differences that might arise due to variable grid spacings among the GCMs. Figure 11 shows  $S_2(r)$  for 2 m air temperatures averaged over this region. Clearly, McICA noise has little impact of fluctuations of surface temperature. The magnitude of differences as a function of McICA noise shown here echo not only other areas of the globe but also resemble closely other surface variables such as precipitation rate and pressure. Note also that all models agree nicely on the magnitude of  $S_2(r)$  and that they all show a well-defined scaling regime between roughly  $r = 1$  h to  $r = 48$  h with  $\zeta(2) \simeq 3/2$  (or  $\beta \simeq 5/2$ ).

The variable whose  $S_2(r)$  exhibited by far the largest sensitivity to McICA noise was, again, low-cloud fraction. Figure 12 shows  $S_2(r)$  for low-cloud fraction averaged for the region bounded by 5°S–5°N and 170°–180°E. This region showed particular sensitivity to McICA noise, as documented by Räisänen *et al.* (2005) for CAM1.8. The most striking aspect of this plot is the hypersensitivity of the GFDL model (regardless of

radiation time-step length) to substantial McICA noise generated by 1COL, where fluctuations on short time-scales are almost an order of magnitude larger than those for REF and  $\zeta(2) \simeq 0$  for all  $r$ , whereas  $\zeta(2) \simeq 1$  for REF for  $r \lesssim 24$  h. Evidently, McICA noise has added a large amount of variance directly to low-cloud fraction estimates coming straight from the cloud parametrization. For its BASIC simulation, however, fluctuations almost match perfectly those of REF for all  $r$ .

The plot in Figure 12 for CAM3 shows that fluctuations in low-cloud fraction for  $r \lesssim 3$  h increase slightly as McICA noise is added. This is perfectly understandable. For larger  $r$ , however, the magnitudes of typical fluctuations decrease as noise increases which seems to be counter-intuitive. As a result, going from REF to 1COL,  $\zeta(2)$  changes from  $\sim 7/6$  to  $\sim 5/6$ . For the CCCma, ECHAM5 and GEM GCMs, the introduction of McICA noise has little effect on  $S_2(r)$ . For ECMWF, however, the magnitude of  $S_2(r)$  is *least* for 1COL for all  $r$ . Again, this seems counter-intuitive for it implies that the addition of McICA noise *suppresses*, albeit slightly, fluctuations of low-cloud fractions. To lesser extents, this is also the case for ECMWF's 2 m temperature as well as low-cloud fraction for CAM3 and GEM.

Another aspect of Figure 12 to note, though not related directly to McICA noise, is that while three of the GCMs (CAM3, CCCma and GEM) exhibit clear diurnal cycles in low-cloud fraction for this region with relative maxima in  $S_2(r)$  at  $r = 12 + 24n$  h and minima at  $r = 24(n+1)$  h, for  $n = 0, 1, 2, 3, \dots$ , the other three lack a diurnal signal. Moreover, the magnitude of  $S_2(r)$ , for any  $r$ , is five to ten times larger for the CCCma and GEM models than for the other models. These differences swamp McICA-generated differences.

To further the points just discussed, Figure 13 shows the global impact of McICA noise on  $\zeta(2)$  for the semi-diurnal range (i.e.  $r \leq 12$  h). Consistent with Figure 12, CAM shows a ubiquitous reduction in  $\zeta(2)$  as noise increases, CCCma, ECHAM5, and GEM show almost no effect, and ECMWF exhibits very little impact as well except a slight, but general, and again counter-intuitive, increase in  $\zeta(2)$  going from REF to 1COL. The odd man out is GFDL where it is now clear that the massive reduction in  $\zeta(2)$  going from REF to 1COL as seen in Figure 12 is actually a global phenomenon that is far from fully recovered by BASIC. Interestingly, however, the corresponding impacts of McICA noise on  $\zeta(2)$  for total precipitation, large-scale precipitation, and 2 m air temperature were all very minor (and not worth showing). Hence, the large changes to fluctuations in low-cloud fraction involved non-precipitating clouds.

## 5. Conclusions

Broadband radiative fluxes computed by the Monte Carlo Independent Column Approximation (McICA) are demonstrably unbiased with respect to the full ICA. As such, McICA represents somewhat of a *dénouement* for the genre of GCM radiative transfer algorithms that aspire

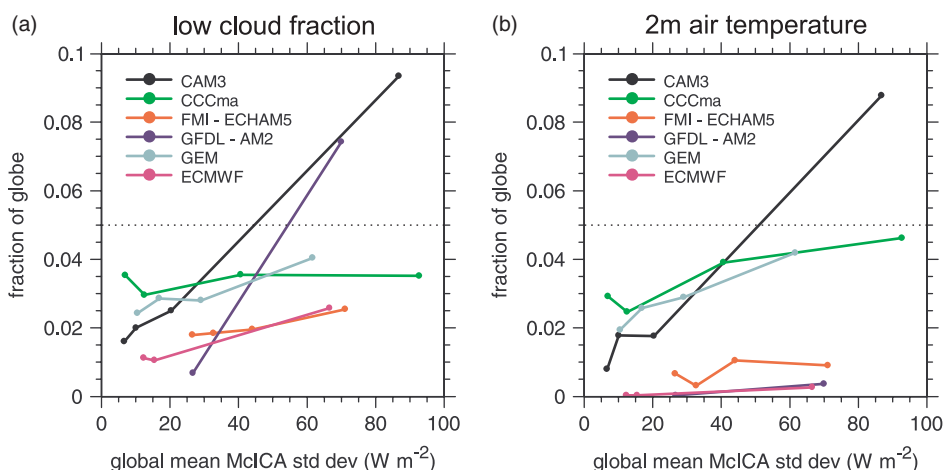


Figure 10. (a) shows fractional areas of the globe exhibiting statistically significant differences for low cloud fraction, at the 95% confidence level, as a function of McICA noise (i.e. standard deviation) associated with net short-wave surface irradiance. (b) is as (a) but for 2 m air temperature. 1COL results have the largest standard deviations. This figure is available in colour online at [www.interscience.wiley.com/journal/qj](http://www.interscience.wiley.com/journal/qj)

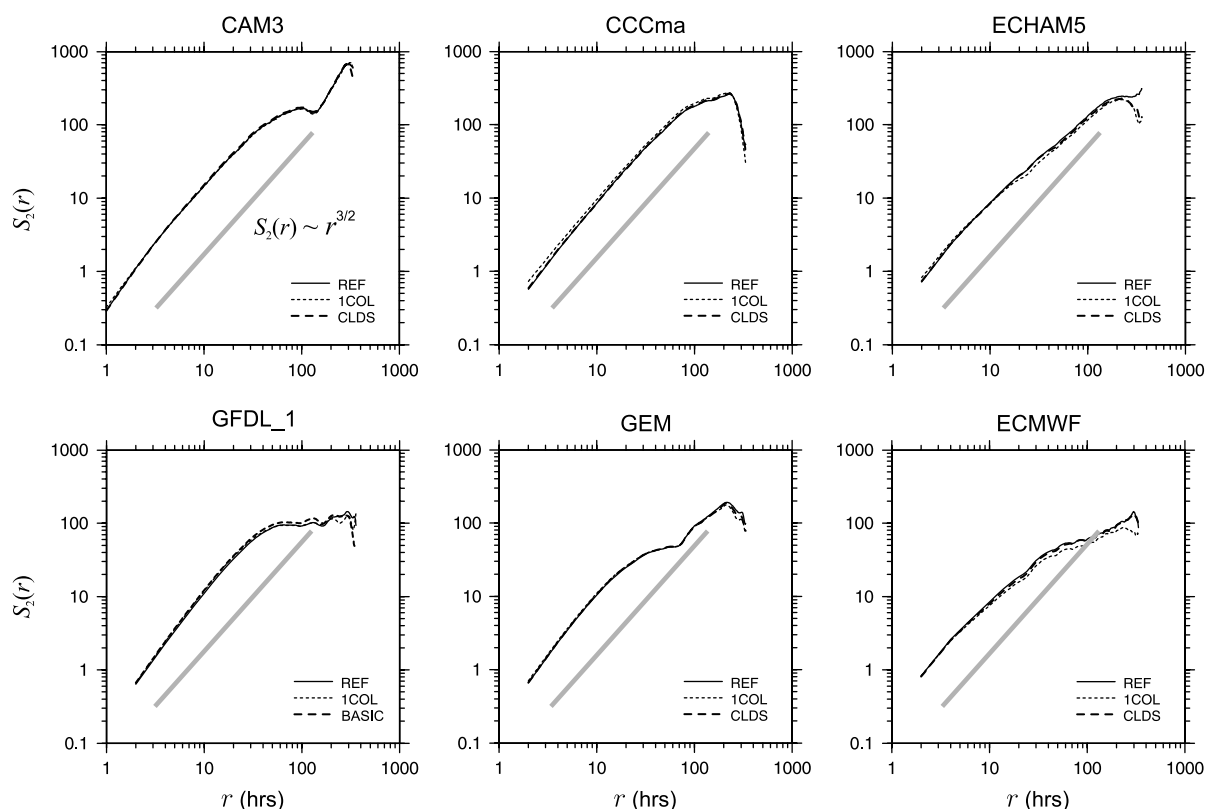


Figure 11. Second-order structure functions for 2 m air temperature averaged over the region 55–65°N, 95–105°W as functions of lag time. Results are shown for three versions of McICA for each GCM. All analyses were for data with 2-hour time steps, except for the CAM3 GCM which had a 1-hour step. The grey line is for reference and highlights the consistency among GCMs and McICA noise levels.

to reproduce full ICA results. The penalty for eliminating biases, however, is random noise. If it so happens that this random noise, extraneous information injected at the GCM's inner-scale, gives rise to statistically significant impacts on GCM simulations, this could render McICA useless. Thus, this paper summarizes the impact of random noise generated by McICA on a diverse group of GCMs.

With the exception of the excessive random noise associated with the 1COL version of McICA, noise

produced by all realizations of McICA, from those most likely to be used in GCMs to those that squelch noise unnecessarily, have no statistically significant affects on short, 14-day, GCM simulations. Even for the 1COL version, not all GCMs showed a significant response to noise and those that did were not impacted significantly for all variables. For instance, the CAM3 GCM showed significant impacts for low cloud and 2 m surface temperature due to noise from 1COL, but its large-scale precipitation was influenced only very weakly. For the BASIC

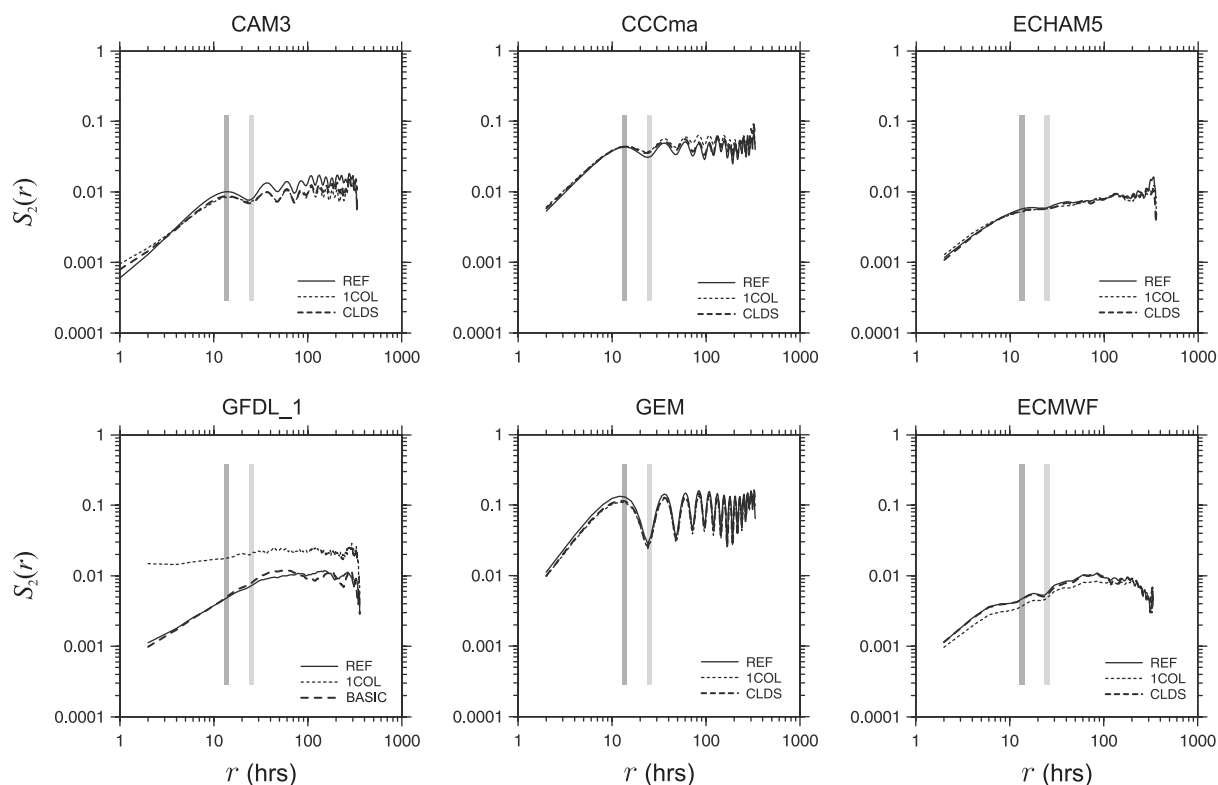


Figure 12. Second-order structure functions for low-cloud fraction averaged over the region  $5^{\circ}\text{S}$ – $5^{\circ}\text{N}$ ,  $170$ – $180^{\circ}\text{E}$  as functions of lag time. Results are shown for three versions of McICA for each GCM. All time series used data with 2-hour time steps, except for the CAM3 GCM which had a 1-hour step. Vertical lines are for reference and indicate periods of 12 and 24 hours.

and CLDS versions of McICA, no model displayed any significant impacts from noise. This corroborates previous studies by Pincus *et al.* (2003; 2006), Räisänen *et al.* (2005), and Morcrette *et al.* (2008) and suggests strongly that, for most applications, McICA-generated noise should be of little or no concern.

There were some peculiar results. For example, the two weather forecast models which operated at the highest spatial resolution actually showed a slight reduction in the magnitude of low-cloud fluctuations at most frequencies as McICA noise increased. This seems counter-intuitive; one would expect introduction of high-frequency noise to increase high-frequency fluctuations in most variables. Another GCM showed only an increase in high-frequency low-cloud fluctuations when noise increased and reductions for frequencies longer than a day.

The experiments performed and assessed here address only the ramifications of noise on fast to intermediate climatic variables. The underlying assumption was that, if these variables were not influenced, it is unlikely that McICA's noise could somehow emerge later on. Of course, given the complexity of modern GCMs, the validity of this assumption cannot be taken fully for granted. Tests should be extended to interactive ocean–atmosphere GCMs integrated over substantial stretches of time. Nevertheless, until it is demonstrated that omission of 3D radiative transfer effects in conventional GCMs is detrimental to simulations (i.e. the ICA is an inadequate standard), results presented here suggest that McICA should be sufficient for most applications of conventional GCMs.

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## References

- Aires F, Rossow WB. 2003. Inferring instantaneous, multivariate and nonlinear sensitivities for the analysis of feedback processes in a dynamical system: The Lorenz model case study. *Q. J. R. Meteorol. Soc.* **129**: 239–275.
- Barker HW. 1996. A parameterization for computing grid-averaged solar fluxes for inhomogeneous marine boundary layer clouds. Part I: Methodology and homogeneous biases. *J. Atmos. Sci.* **53**: 2289–2303.
- Barker HW, Stephens GL, Fu Q. 1999. The sensitivity of domain-averaged solar fluxes to assumptions about cloud geometry. *Q. J. R. Meteorol. Soc.* **125**: 2127–2152.
- Barker HW, Pincus R, Morcrette J-J. 2002. 'The Monte Carlo Independent Column Approximation: Application within large-scale models'. In proceedings of the GCSS-ARM Workshop on the Representation of Cloud Systems in Large-Scale Models, May 2002, Kananaskis, AB, Canada. Available at: <http://www.met.utah.edu/skrueger/gcss-2002/Extended-Abstracts.pdf>.
- Barker HW, Stephens GL, Partain PT, Bergman JW, Bonnel B, Campana K, Clothiaux EE, Clough S, Cusack S, Delamere J, Edwards J, Evans KF, Fouquart Y, Freidenreich S, Galin V, Hou Y,



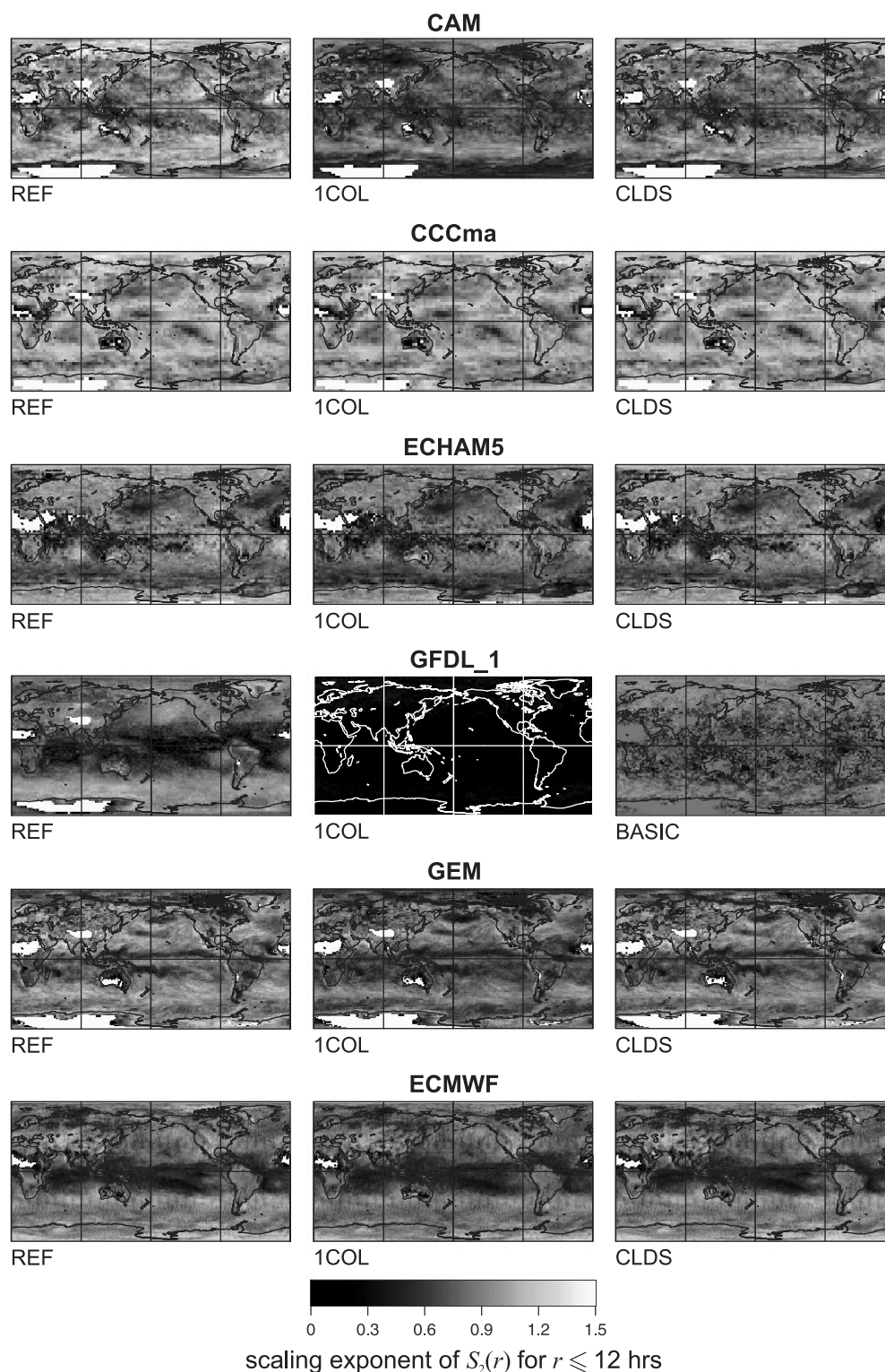


Figure 13. Global plots of the scaling exponent  $\zeta(2)$  in Equation (9) for low-cloud fraction for three versions of McICA. These quantities correspond to the scaling regions shown in Figure 12 for  $r \leq 12$  h.

Kato S, Li J, Mlawer E, Morcrette J-J, O'Hirok W, Räisänen P, Ramaswamy V, Ritter B, Rozanov E, Schlesinger M, Shibata K, Sporyshev P, Sun Z, Wendisch M, Wood N, Yang F. 2003. Assessing 1D atmospheric solar radiative transfer models: Interpretation and handling of unresolved clouds. *J. Climate* **16**: 2676–2699.

Cahalan RF, Ridgway W, Wiscombe WJ, Gollmer S, Harshvardhan. 1994. Independent pixel and Monte Carlo estimates of stratocumulus albedo. *J. Atmos. Sci.* **51**: 3776–3790.

Cairns B, Lacis AA, Carlson BE. 2000. Absorption within inhomogeneous clouds and its parameterization in general circulation models. *J. Atmos. Sci.* **57**: 700–714.

Cole JNS, Barker HW, O'Hirok W, Clothiaux EE, Khairoutdinov MF, Randall DA. 2005. Atmospheric radiative transfer through global arrays of 2D clouds. *Geophys. Res. Lett.* **32**: L19817, doi: 10.1029/2005GL023329.

- Collins WD. 2001. Parameterization of generalized cloud overlap for radiative calculations in general circulation models. *J. Atmos. Sci.* **58**: 3224–3242.
- Davis A, Gabriel P, Lovejoy S, Schertzer D, Austin G. 1990. Discrete angle radiative transfer - Part III: Numerical results and applications. *J. Geophys. Res.* **95**: 11729–11742.
- Fu Q, Liou K-N. 1992. On the correlated *k*-distribution method for radiative transfer in inhomogeneous atmospheres. *J. Atmos. Sci.* **49**: 2139–2156.
- Geleyn J-F, Hollingsworth A. 1979. An economical analytical method for the computation of the interaction between scattering and line absorption of radiation. *Contrib. Atmos. Phys.* **52**: 1–16.
- Intergovernmental Panel on Climate Change (IPCC). 2007. Working Group I Report (WG1): Climate Change 2007: The Physical Science Basis. Available at <http://ipcc-wg1.ucar.edu/wg1/wg1-report.html>.
- Marchuk GI, Mikhailov GA, Nazaraev MA, Darbinjan RA, Kargin BA, Elepov BS. 1980. *Monte Carlo methods in atmospheric optics*. Ser. Opt. Sci. **12**: Springer-Verlag: Berlin.
- Marshak A, Davis AB. 2005. Scale-by-scale analysis and fractal cloud models. Pp. 653–664 in *3D Radiative Transfer in Cloudy Atmospheres*, Marshak A, Davis AB (eds.) Springer: Heidelberg.
- Meador WE, Weaver WR. 1980. Two-stream approximations to radiative transfer in planetary atmospheres: A unified description of existing methods and a new improvement. *J. Atmos. Sci.* **37**: 630–643.
- Morcrette J-J. 2000. On the effects of the temporal and spatial sampling of radiation fields on the ECMWF forecasts and analyses. *Mon. Weather Rev.* **128**: 876–887.
- Morcrette J-J, Fouquart Y. 1986. The overlapping of cloud layers in shortwave radiation parameterizations. *J. Atmos. Sci.* **43**: 321–328.
- Morcrette J-J, Barker HW, Cole JNS, Iacono MJ, Pincus R. 2008. Impact of the new radiation package, McRAD, in the ECMWF integrated forecast system. *Mon. Weather Rev.* In press.
- Oreopoulos L, Barker HW. 1999. Accounting for subgrid-scale cloud variability in a multi-layer, 1D solar radiative transfer algorithm. *Q. J. R. Meteorol. Soc.* **125**: 301–330.
- Phillips TJ, Potter GL, Williamson DL, Cederwall RT, Boyle JS, Fiorino M, Hnilo JJ, Olson JG, Xie S, Yio JJ. 2004. 'The CCPP-ARM Parameterization Testbed (CAPT): Evaluating climate models in a weather forecasting framework'. 15th Symposium on Global Change and Climate Variations. <http://ams.confex.com/ams/pdfpapers/73412.pdf>
- Pincus R, Barker HW, Morcrette J-J. 2003. A fast, flexible, approximate technique of computing radiative transfer for inhomogeneous clouds. *J. Geophys. Res.* **108**: 4376, doi:10.1029/2002JD003322.
- Pincus R, Hemler R, Klein SA. 2006. Using stochastically generated subcolumns to represent cloud structure in a large-scale model. *Mon. Weather Rev.* **134**: 3644–3656.
- Randall D, Khairoutdinov M, Arakawa A, Grabowski W. 2003. Breaking the cloud parameterization deadlock. *Bull. Am. Meteorol. Soc.* **84**: 1547–1564.
- Räisänen P, Barker HW, Khairoutdinov M, Li J, Randall DA. 2004. Stochastic generation of subgrid-scale cloudy columns for large-scale models. *Q. J. R. Meteorol. Soc.* **130**: 2047–2067.
- Räisänen P, Barker HW. 2004. Evaluation and optimization of sampling errors for the Monte Carlo independent column approximation. *Q. J. R. Meteorol. Soc.* **130**: 2069–2086.
- Räisänen P, Barker HW, Cole J. 2005. The Monte Carlo Independent Column Approximation's conditional random noise: Impact on simulated climate. *J. Climate* **17**: 4715–4730.
- Stephens GL. 1988. Radiative transfer through arbitrary shaped optical media: Part II: Group theory and simple closures. *J. Atmos. Sci.* **45**: 1837–1848.
- Stephens GL. 2005. Cloud feedbacks in the climate system: A critical review. *J. Climate* **18**: 237–273.
- Stephens GL, Gabriel PM, Tsay S-C. 1991. Statistical radiative transfer in one-dimensional media and its application to the terrestrial atmosphere. *Trans. Theory Statist. Phys.* **20**: 139–175.
- Stubenrauch CJ, Del Genio AD, Rossow WB. 1997. Implementation of subgrid cloud vertical structure inside a GCM and its effect on the radiation budget. *J. Climate* **10**: 273–287.
- Tian L, Curry JA. 1989. Cloud overlap statistics. *J. Geophys. Res.* **94**: 9925–9935.
- von Storch H, Zwiers FW. 1999. *Statistical Analysis in Climate Research*. Cambridge University Press: Cambridge, UK.
- Wiscombe WJ. 1977. 'The delta-Eddington approximation for a vertically inhomogeneous atmosphere'. Tech. Note TN-121+STR. NCAR: Boulder, USA.